

Hadronic molecules with a short-range force by a quark model

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in collaboration with

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Social Work, ⁵Showa Pharmaceutical U.

Y.Y, A. Giachino, A. Hosaka, E. Santopinto, S. Takeuchi, M. Takizawa,
Phys.Rev. D96 (2017), 114031

Hadrons and Nuclear Physics meet ultracold atoms: a French
Japanese workshop

Hadronic molecules + Compact state

1 Introduction

- Exotic hadron
- Hidden-charm pentaquark

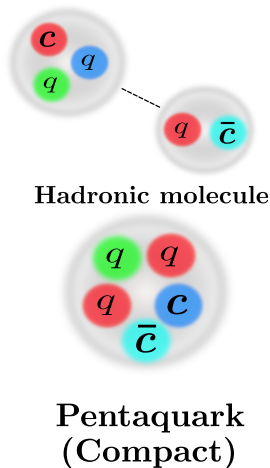
2 Model setup

- Heavy Quark Spin Symmetry and OPEP
- Compact 5-quark potential

3 Numerical results

- Hidden-charm molecules

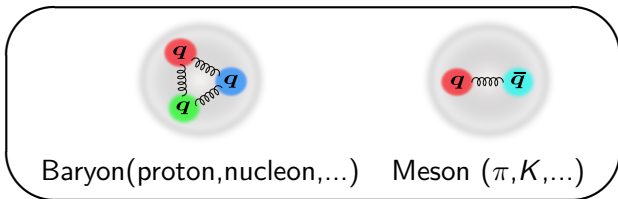
4 Summary



Conventional and Exotic hadrons

Introduction: Exotic hadron

- Hadron: Composite particle of **Quarks** and **Gluons**
- Constituent quark model (Baryon(qqq) and Meson $q\bar{q}$) has been successfully applied to the hadron spectra!

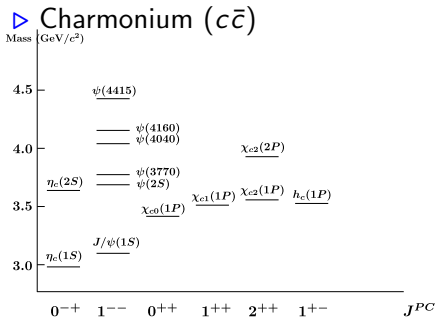


- ▶ Quark potential (Coulomb type + Linear potential)

$$V_q(r) = -\frac{a}{r} + br + c + d \frac{\vec{S}_1 \cdot \vec{S}_2}{m_1 m_2} \delta^{(3)}(r) + \dots$$

Exotic hadrons in the heavy quark region

Introduction: Exotic hadron



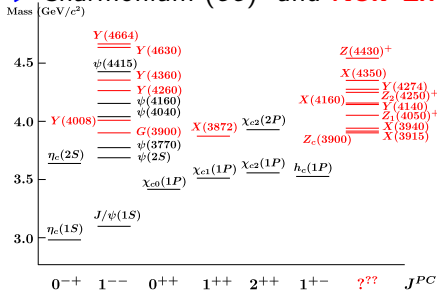
Charmonium $c\bar{c}$

N. Brambilla, *et al.* Eur.Phys.J.C **71**(2011)1534
S. Godfrey and N. Isgur, PRD**32**(1985)189

Exotic hadrons in the heavy quark region

Introduction: Exotic hadron

Charmonium ($c\bar{c}$) and **New Exotic hadrons X, Y, Z**



Charmonium $c\bar{c}$
and
Exotic hadrons ($\neq c\bar{c}$)

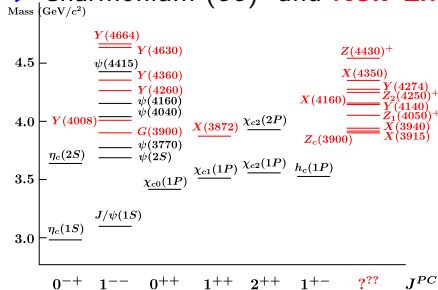
X, Y, Z

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Exotic hadrons in the heavy quark region

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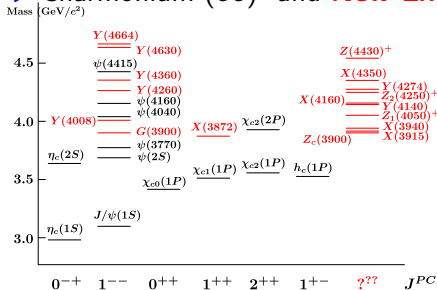
► **What is the structure of exotic hadrons?**

⇒ Multiquark states?

Exotic hadrons in the heavy quark region

Introduction: Exotic hadron

Charmonium ($c\bar{c}$) and **New Exotic hadrons X, Y, Z**

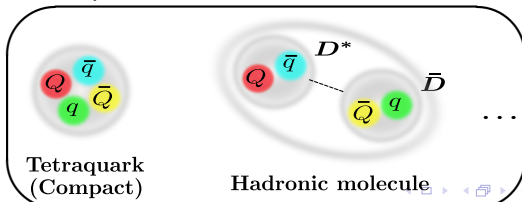


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► What is the structure of exotic hadrons?

⇒ Multiquark states?



Observation of two hidden-charm pentaquarks !!

Introduction: pentaquark

PRL 115, 072001 (2015)

PHYSICAL REVIEW LETTERS

week ending
14 AUGUST 2015



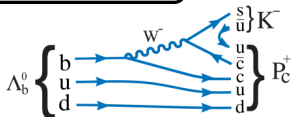
Observation of J/ψ Resonances Consistent with Pentaquark States in $\Lambda_b^0 \rightarrow J/\psi K^- p$ Decays

R. Aaij *et al.**

(LHCb Collaboration)

(Received 13 July 2015; published 12 August 2015)

$\Lambda_b^0 \rightarrow K^- P_c^+$ decay



Two ($c\bar{u}ud$) Pentaquarks!

$P_c(4380)$: $M=4380 \pm 8 \pm 29$ MeV
 $\Gamma = 205 \pm 18 \pm 86$ MeV

$P_c(4450)$: $M=4449.8 \pm 1.7 \pm 2.5$ MeV
 $\Gamma = 39 \pm 5 \pm 19$ MeV

J^P : $(\frac{3}{2}^-, \frac{5}{2}^+)$, $(\frac{3}{2}^+, \frac{5}{2}^-)$ or $(\frac{5}{2}^+, \frac{3}{2}^-)$ *Opposite parity

best fit!

- $P_c(4380)$ and $P_c(4450)$ obtained near $\underline{\bar{D}\Sigma_c^*}$ and $\underline{\bar{D}^*\Sigma_c}$

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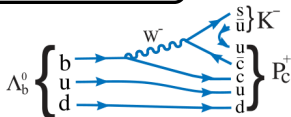
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best fit!

- $P_c(4380)$ and $P_c(4450)$ obtained near $\bar{D}\Sigma_c^*$ and $\bar{D}^*\Sigma_c$
- ▶ Possible existence of Exotic baryons in the hidden-charm sector?

Theoretical discussions of the hidden-charm baryons

Introduction: pentaquark

Proposals of various structures!

H.X.Chen, *et al.*, Phys.Rept. **639**(2016)1, A.Esposito, *et al.*, Phys.Rept. **668**(2016)1, A.Ali, *et al.*, PPNP **97**(2017)123

- Compact pentaquark ($c\bar{c}qqq$)?

S.G.Yuan, *et al.* (2012), L.Maiani, *et al.* (2015),

S.Takeuchi, *et al.* (2017), J. Wu, *et al.* (2017),

...

- Hadronic molecule ($\bar{D}\Sigma_c^*$, $\bar{D}^*\Sigma_c, \dots$)?

J.-J.Wu *et al.*, (2010) (2011), C. Garcia-Recio, *et al.* (2013),

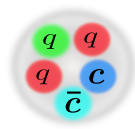
R. Chen, *et al.* (2015), Y.Shimizu, *et al.* (2016) (2017),

...

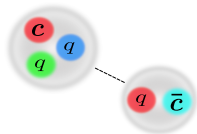
- Kinematical effect? Cusp?
(Non-resonant explanation)

F.K.Guo, *et al.* (2015), X.H.Liu, *et al.* (2016),

...



Pentaquark
(Compact)



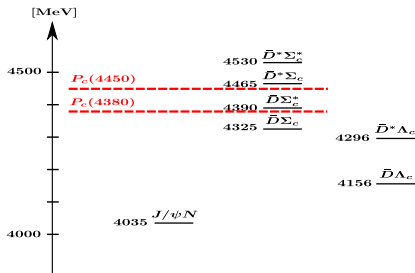
Hadronic molecule

Exotic states near thresholds \rightarrow Molecules?

Introduction: pentaquark

\triangleright e.g. $P_c(4380)$, $P_c(4450)$

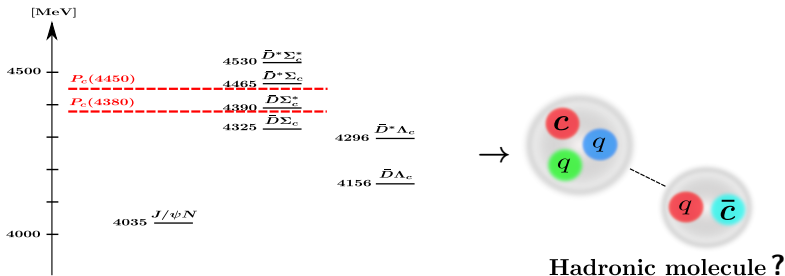
\rightarrow close to **the meson-baryon thresholds**



Exotic states near thresholds \rightarrow Molecules?

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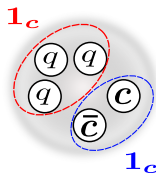
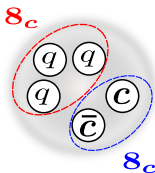
- Exotic state may be a loosely bound state of the meson-baryon.
 - \Rightarrow Analogous to atomic nuclei (Deuteron: $B \sim 2.2$ MeV)

Importance of Hadron-hadron interaction (not known...)

Compact state: 5-quark configuration

Introduction: pentaquark

- S. Takeuchi and M. Takizawa, PLB**764** (2017) 254-259.
 P_c states by the quark cluster model
- 5-quark configurations

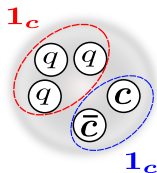
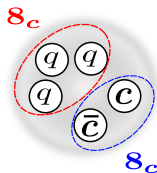


$$S_{q^3} = 1/2, 3/2, S_{c\bar{c}} = 0, 1 \quad S_{q^3} = 1/2, S_{c\bar{c}} = 0, 1$$

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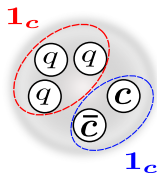
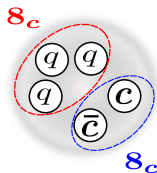
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- $[q^3 8_c 3/2]$: Color magnetic int. is attractive!

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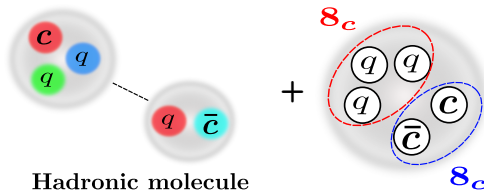
- $[q^3 8_c 3/2]$: Color magnetic int. is attractive!
 \Rightarrow Couplings to (qqc) baryon- $(q\bar{c})$ meson, e.g. $\bar{D}\Sigma_c$, are allowed!

Mixing of Compact state and Hadronic Molecule!

Model setup in this study

- Hadronic molecule (MB) + Compact state ($5q$)

$MB + 5q$

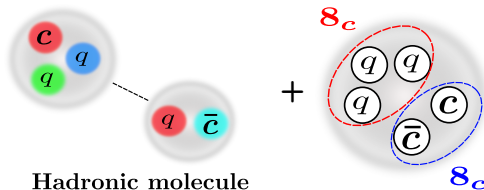


Hadronic molecule

Model setup in this study

- **Hadronic molecule (MB) + Compact state ($5q$)**
⇒ MB coupled to $5q$ (Feshbach Projection)

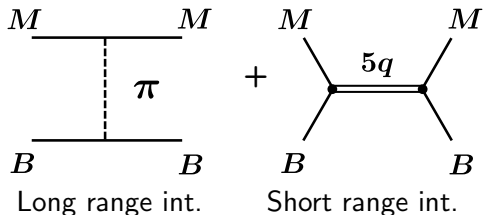
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Interaction of hadrons (M and B)

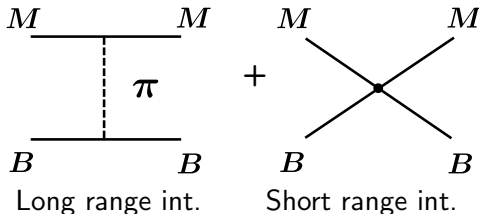


- ▶ **Long range** interaction: One pion exchange potential (OPEP)
- ▶ **Short range** interaction: $5q$ potential

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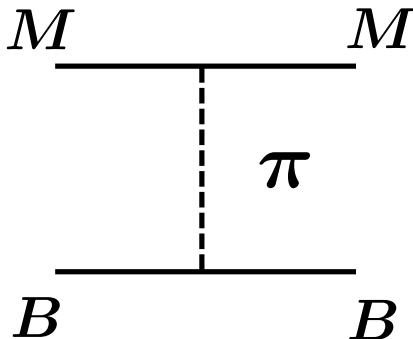
Interaction of hadrons (M and B)



- ▶ **Long range** interaction: One pion exchange potential (OPEP)
- ▶ **Short range** interaction: $5q$ potential (→ **Local Gaussian**)
(* Other int. (**double counting...**) → [Future work](#))

MB bound states: Role of the $5q$ potential (Spin structure)

1. Long range force: One pion exchange potential



- Exchanging light meson π ($m_\pi \sim 140$ MeV)
- Driving force to bind Atomic nucleus

Heavy Quark Spin Symmetry

Heavy Quark Spin Symmetry

Charm (c), Bottom (b), Top (t)

Heavy Quark Spin Symmetry

Charm (c), Bottom (b), Top (t)



1. Coupled channels of MB
2. Tensor force (OPEP)

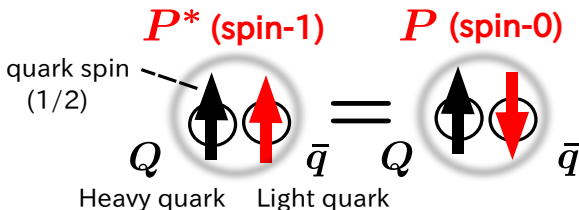
Heavy Quark Spin Symmetry and Mass degeneracy

HQS and OPEP

Heavy Quark Spin Symmetry (HQS)

N.Isgur, M.B.Wise, PLB232(1989)113

- **Suppression of Spin-spin force** in $m_Q \rightarrow \infty$.
 \Rightarrow **Mass degeneracy** of hadrons with the different J
- e.g. $Q\bar{q}$ meson



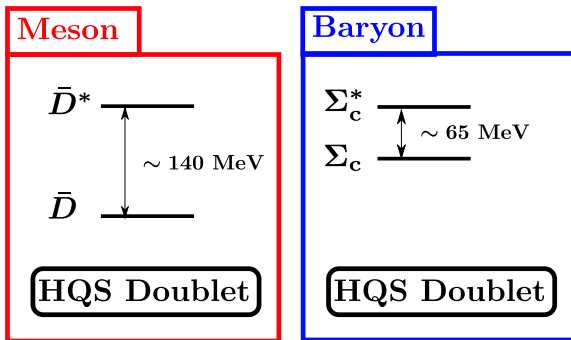
\Rightarrow Mass degeneracy of spin-0 and spin-1 states!

- Charm sector: $\bar{D}(0^-) - \bar{D}^*(1^-)$, $\Sigma_c(1/2^+) - \Sigma_c^*(3/2^+)$

Mass degeneracy $\rightarrow \bar{D} - \bar{D}^*$, $\Sigma_c - \Sigma_c^*$ mixing!

HQS and OPEP

- $\bar{D} - \bar{D}^*$ and $\Sigma_c - \Sigma_c^*$ mixing in the $\bar{D}Y_c$ system

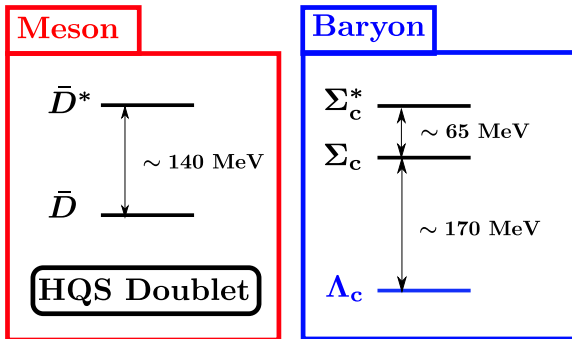


- Coupled channels of $\bar{D}\Sigma_c$, $\bar{D}\Sigma_c^*$, $\bar{D}^*\Sigma_c$ and $\bar{D}^*\Sigma_c^*$!

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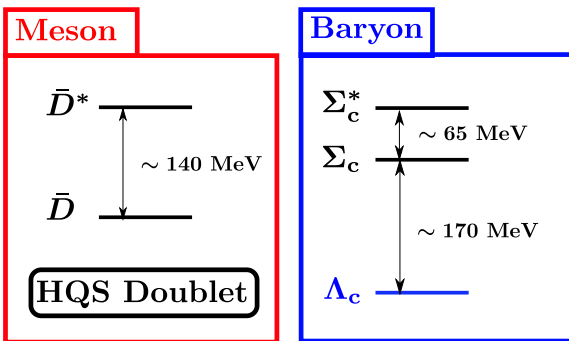


- Coupled channels of $\bar{D}\Sigma_c$, $\bar{D}\Sigma_c^*$, $\bar{D}^*\Sigma_c$ and $\bar{D}^*\Sigma_c^*$!
- In addition, Λ_c (cqq): $\bar{D}^{(*)}\Lambda_c$ channel!?

Mass degeneracy $\rightarrow \bar{D} - \bar{D}^*$, $\Sigma_c - \Sigma_c^*$ mixing!

HQS and OPEP

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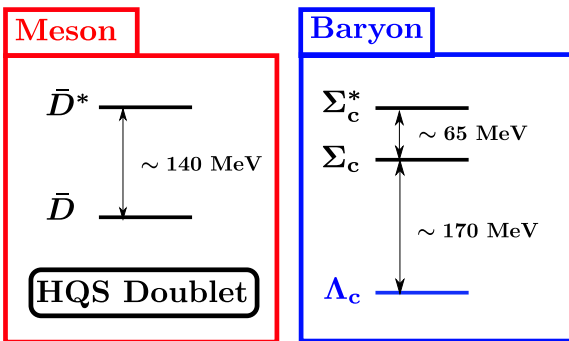
- ▷ 6 meson-baryon components

(1) $\bar{D}\Lambda_c$, (2) $\bar{D}^*\Lambda_c$, (3) $\bar{D}\Sigma_c$, (4) $\bar{D}\Sigma_c^*$,
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HQS and OPEP

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(5) $\bar{D}^*\Sigma_c$, (6) $\bar{D}^*\Sigma_c^* \rightarrow$ Coupled by OPEP!

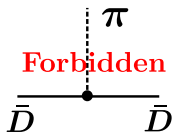
Heavy hadron- π coupling

HQS and OPEP

- Effective Lagrangians: Heavy hadron and π

R. Casalbuoni *et al.*, Phys.Rept.**281** (1997)145, T. M. Yan, *et al.*, PRD**46**(1992)1148

Y.-R.Liu and M.Oka, PRD**85**(2012)014015



- ▶ Heavy meson: $\bar{D}^{(*)}\bar{D}^{(*)}\pi$ ($DD\pi$: Parity violation)

$$\mathcal{L}_{\pi HH} = -\frac{g_{\pi}}{2f_{\pi}} \text{Tr} [H\gamma_{\mu}\gamma_5\partial^{\mu}\hat{\pi}\bar{H}], \quad H = \frac{1+\not{\epsilon}}{2} [D_{\mu}^{*}\gamma^{\mu} - D\gamma_5]$$

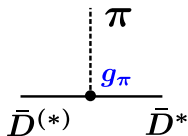
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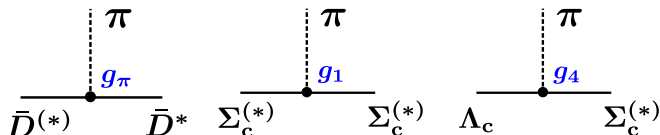
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- Heavy baryon: $\Sigma_c^{(*)}\Sigma_c^{(*)}\pi, \Lambda_c\Sigma_c^{(*)}\pi$ ($\Lambda_c\Lambda_c\pi$: Isospin breaking)

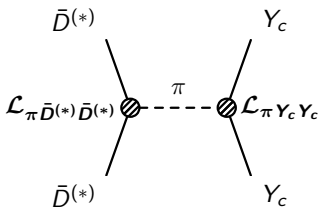
$$\mathcal{L}_{\pi BB} = -\frac{3}{4f_\pi} \mathbf{g}_1 (i v_\kappa) \varepsilon^{\mu\nu\lambda\kappa} \text{tr} [\bar{S}_\mu \partial_\nu \hat{\pi} S_\lambda] - \frac{g_4}{2f_\pi} \text{tr} [\bar{S}^\mu \partial_\mu \hat{\pi} \Lambda_c] + \text{H.c.},$$

$$\mathbf{S}_\mu = \Sigma_{c\mu}^* - \frac{1}{\sqrt{3}} (\gamma_\mu + \mathbf{v}_\mu) \gamma_5 \Sigma_c, \quad g_\pi = 0.59, g_1 = 1.00, g_4 = 1.06$$

$\bar{D}^{(*)} Y_c$ Interaction: Long range force

HQS and OPEP

- One pion exchange potential



$\bar{D}^{(*)}$: \bar{D} or \bar{D}^*

Y_c : Λ_c , Σ_c or Σ_c^*

$$V_{\bar{D}^{(*)} Y_c - \bar{D}^{(*)} Y_c}^{\pi} = G \left[\vec{S}_1 \cdot \vec{S}_2 C(r) + S_{S_1 S_2} T(r) \right]$$

(Contact term is removed)

- Form factor with Cutoff Λ (determined by the hadron size)

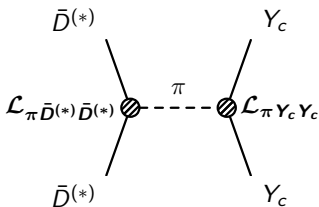
$$F(q^2) = \frac{\Lambda^2 - m_{\pi}^2}{\Lambda^2 - q^2}, \quad \Lambda_{\bar{D}} \sim 1130 \text{ MeV}, \Lambda_{Y_c} \sim 840 \text{ MeV}$$

Y.Y, A. Giachino, A. Hosaka, E. Santopinto, S. Takeuchi, M. Takizawa, Phys.Rev. **D96** (2017), 114031

$\bar{D}^{(*)} Y_c$ Interaction: Long range force

HQS and OPEP

- One pion exchange potential **with Tensor force!**



$\bar{D}^{(*)}$: \bar{D} or \bar{D}^*

Y_c : Λ_c , Σ_c or Σ_c^*

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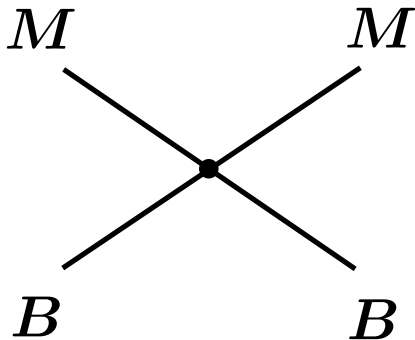
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$$F(q^2) = \frac{\Lambda^2 - m_{\pi}^2}{\Lambda^2 - q^2}, \quad \Lambda_{\bar{D}} \sim 1130 \text{ MeV}, \Lambda_{Y_c} \sim 840 \text{ MeV}$$

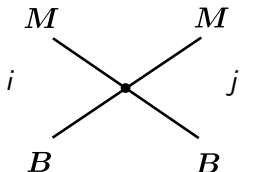
Y.Y, A. Giachino, A. Hosaka, E. Santopinto, S. Takeuchi, M. Takizawa, Phys.Rev. **D96** (2017), 114031

2. Short range force: 5-quark potential



Model: 5-quark potential

- 5-quark potential \Rightarrow **Local Gaussian potential** is employed.
- ▷ Massive M_{5q} (few hundred MeV above $\bar{D}^*\Sigma_c^*$) \rightarrow **Attractive**


$$\Rightarrow -f S_i S_j e^{-\alpha r^2}$$

Channel $i, j = \bar{D}^{(*)}\Lambda_c, \bar{D}^{(*)}\Sigma_c^{(*)}$ with S -wave

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$$\begin{array}{ccc} M & & M \\ & \diagdown & / \\ & \bullet & \\ & / & \diagdown \\ B & & B \end{array} \quad i \quad j \quad \Rightarrow -f S_i S_j e^{-\alpha r^2}$$

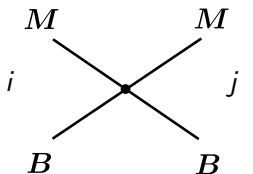
Channel $i, j = \bar{D}^{(*)}\Lambda_c, \bar{D}^{(*)}\Sigma_c^{(*)}$ with S -wave

Free Parameters

Strength f and Gaussian para. α (\rightarrow may be fixed in the future)
(f vs E will be shown latter. $\alpha = 1 \text{ fm}^{-2}$ is fixed.)

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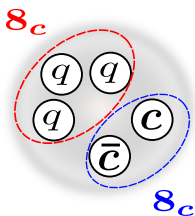
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Spectroscopic factors S_i (Spin structure)

$5q$ potential

- Spin of $5q$ states $\rightarrow S_{c\bar{c}}$ and S_{3q} configuration
e.g. for $J^P = 1/2^-$, (i), (ii), (iii)



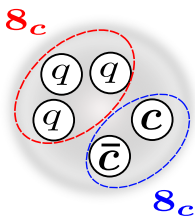
$$J^P = 1/2^-$$

	$S_{c\bar{c}}$	S_{3q}
type (i)	0	1/2
(ii)	1	1/2
(iii)	1	3/2

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- Overlap** of the spin wavefunctions of 5-quark state and $\bar{D}Y_c$

$$S_i = \langle (\bar{D}Y_c)_i | 5q \rangle$$

\Rightarrow Relative strength of couplings to $\bar{D}Y_c$ channel

Spectroscopic factor S_i (Spin structure)

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- S-factor = Relative strength to $\bar{D}^{(*)}\Lambda_c$ and $\bar{D}^{(*)}\Sigma_c^{(*)}$

Table: Spectroscopic factors S_i for each meson-baryon channel.

J		$S_{c\bar{c}}$	S_{3q}	$\bar{D}\Lambda_c$	$\bar{D}^*\Lambda_c$	$\bar{D}\Sigma_c$	$\bar{D}\Sigma_c^*$	$\bar{D}^*\Sigma_c$	$\bar{D}^*\Sigma_c^*$
1/2	(i)	0	1/2	0.4	0.6	-0.4	—	0.2	-0.6
	(ii)	1	1/2	0.6	-0.4	0.2	—	-0.6	-0.3
	(iii)	1	3/2	0.0	0.0	-0.8	—	-0.5	0.3
3/2	(i)	0	3/2	—	0.0	—	-0.5	0.6	-0.7
	(ii)	1	1/2	—	0.7	—	0.4	-0.2	-0.5
	(iii)	1	3/2	—	0.0	—	-0.7	-0.8	-0.2
5/2	(i)	1	3/2	—	—	—	—	—	-1.0

Spectroscopic factor S_i (Spin structure)

$5q$ potential

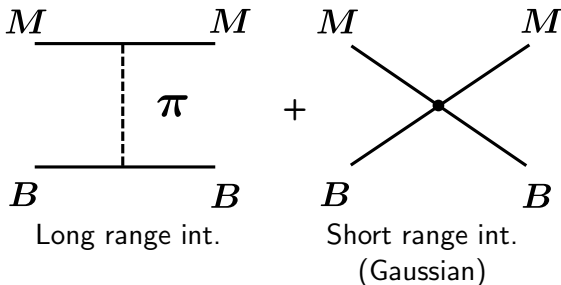
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- $\bar{D}Y_c$ with **Large S_i** will play an important role.

Numerical Results for Hidden-charm sector



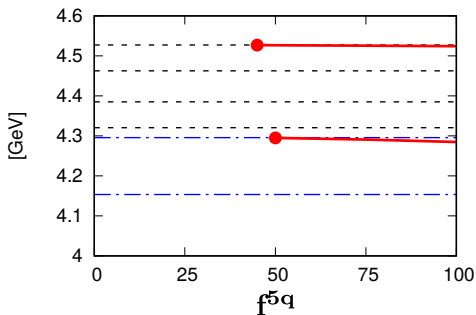
Bound state and Resonance

- Coupled-channel Schrödinger equation for $\bar{D}\Lambda_c$, $\bar{D}^*\Lambda_c$, $\bar{D}\Sigma_c$, $\bar{D}\Sigma_c^*$, $\bar{D}^*\Sigma_c$, $\bar{D}^*\Sigma_c^*$ (6 MB components).
- For $J^P = 1/2^-, 3/2^-, 5/2^-$ (Negative parity)

Results (f^{5q} vs E) of charm $\bar{D}Y_c$ for $J^P = 1/2^-$

- π exchange + V^{5q} [type (i)]

$$(i) (S_{c\bar{c}}, S_{3q}) = (0, \frac{1}{2})$$

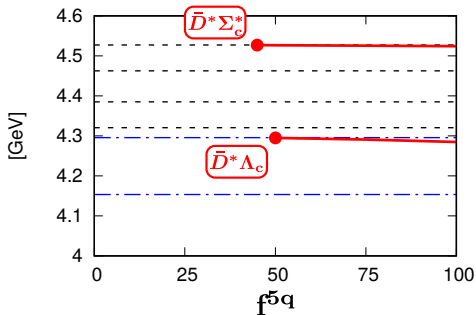


- No state in small f^{5q}
- π exchange is not enough to produce a bound state
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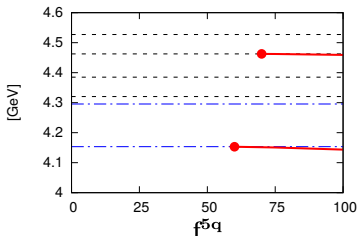


- No state in small f^{5q}
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⇔ **Large S-factor** (Spin structure)

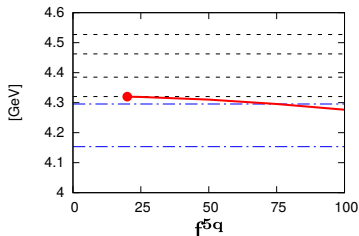
Results (f^{5q} vs E) of charm $\bar{D}Y_c$ for $J^P = 1/2^-$

- π exchange + V^{5q} [types (ii), (iii)]

(ii) $(S_{c\bar{c}}, S_{3q}) = (1, \frac{1}{2})$



(iii) $(S_{c\bar{c}}, S_{3q}) = (1, \frac{3}{2})$

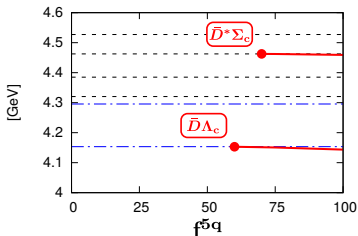


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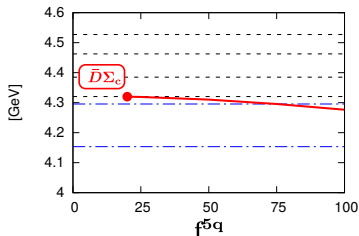
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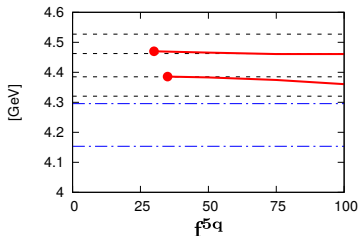


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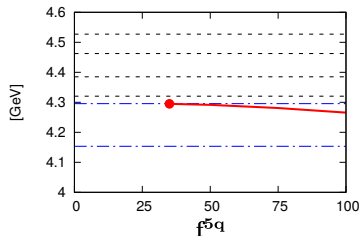
Results (f^{5q} vs E) of charm $\bar{D}Y_c$ for $J^P = 3/2^-$

- π exchange + V^{5q} (i), (ii), (iii)

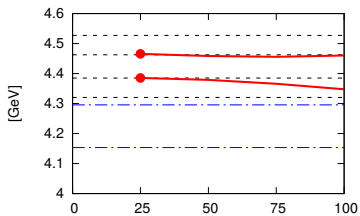
(i) $(S_{c\bar{c}}, S_{3q}) = (0, \frac{3}{2})$



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(iii) $(S_{c\bar{c}}, S_{3q}) = (1, \frac{3}{2})$

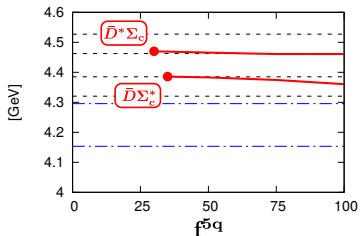


- No state in small f^{5q}
- ⇒ States appear near the thresholds

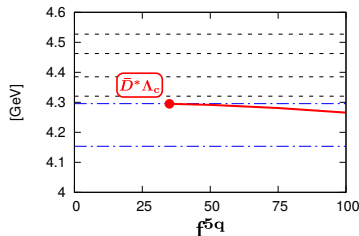
Results (f^{5q} vs E) of charm $\bar{D}Y_c$ for $J^P = 3/2^-$

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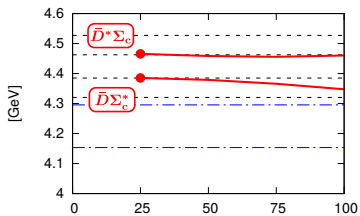
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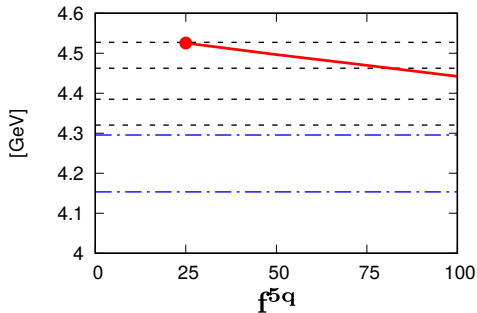


- \bullet No state in small f^{5q}
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- \Leftrightarrow **Large S-factor**

Results (f^{5q} vs E) of charm $\bar{D}Y_c$ for $J^P = 5/2^-$

- Charm $\bar{D}Y_c$ for $J^P = 5/2^-$, One $5q$ state

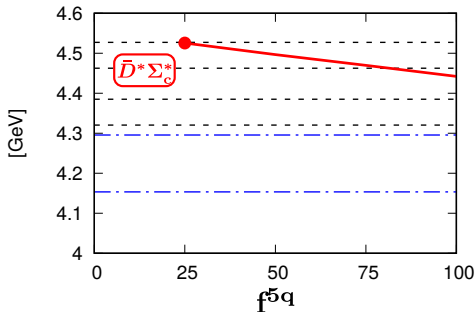
$$(S_{c\bar{c}}, S_{3q}) = (1, \frac{3}{2})$$



Results (f^{5q} vs E) of charm $\bar{D}Y_c$ for $J^P = 5/2^-$

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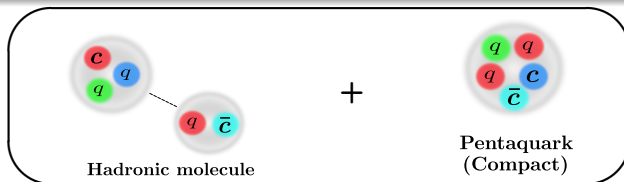
$$(S_{c\bar{c}}, S_{3q}) = (1, \frac{3}{2})$$



Summary of the hidden-charm sector

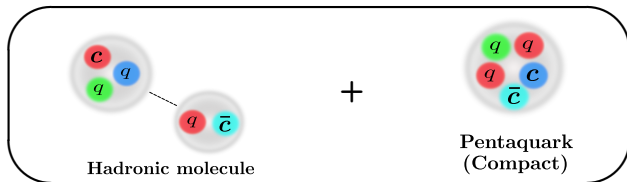
- OPEP is not strong enough to produce a state.
- The importance of the $5q$ potential
 \Rightarrow States below the MB thresholds \leftarrow **large S-factor**

Summary



- **Hadron interaction** is important in Hadronic molecules.
- Coupling to Compact $5q$ state is introduced as the short range interaction.
- Introducing **6 meson-baryon components**:
Multiplet of the HQS, $\bar{D}\Sigma_c, \bar{D}\Sigma_c^*, \bar{D}^*\Sigma_c, \bar{D}^*\Sigma_c^* + \bar{D}\Lambda_c, \bar{D}^*\Lambda_c$
- Interaction: **π exchange** as a long range int., and **the compact 5-quark potential** as a short range int.
- For the hidden-charm, the π exchange is not enough to produce the states. **Importance of the $5q$ potential (Spin structure)**.

Summary



- Future Works

- ▶ Treatment of the $5q$ potential

1. Determining the strength f^{5q} (Quark model?)

2. Energy dependent $5q$ potential

3. Including the $J/\psi N$ channel

- ▶ Other short range interaction (double counting)

Y. Yamaguchi, A. Giachino, A. Hosaka, E. Santopinto, S. Takeuchi, M. Takizawa,
Phys.Rev. D**96** (2017), 114031

Thank you for your kind attention.

Back up

Coupled-channels

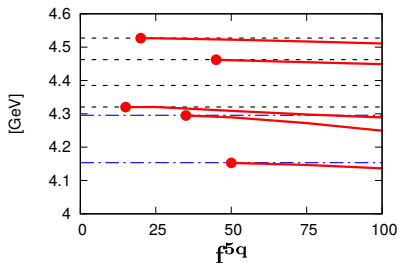
Channels	$\bar{D}Y_c(2S+1L)$	
$1/2^-$	$\bar{D}\Lambda_c(2S), \bar{D}^*\Lambda_c(2S),$ $\bar{D}\Sigma_c(2S), \bar{D}\Sigma_c^*(4D),$ $\bar{D}^*\Sigma_c(2S, 4D), \bar{D}^*\Sigma_c^*(2S, 4D, 6D)$	(10 ch)
$3/2^-$	$\bar{D}\Lambda_c(2D), \bar{D}^*\Lambda_c(4S, 2D, 4D),$ $\bar{D}\Sigma_c(2D), \bar{D}\Sigma_c^*(4S, 4D),$ $\bar{D}^*\Sigma_c(4S, 2D, 4D), \bar{D}^*\Sigma_c^*(4S, 2D, 4D, 6D, 6G)$	(15 ch)
$5/2^-$	$\bar{D}\Lambda_c(2D), \bar{D}^*\Lambda_c(2D, 4D, 4G),$ $\bar{D}\Sigma_c(2D), \bar{D}\Sigma_c^*(4D, 4G),$ $\bar{D}^*\Sigma_c(2D, 4D, 4G), \bar{D}^*\Sigma_c^*(6S, 2D, 4D, 6D, 4G, 6G)$	(16 ch)

- 6 $\bar{D}Y_c$ channels: $\bar{D}\Lambda_c, \bar{D}^*\Lambda_c, \bar{D}\Sigma_c, \bar{D}\Sigma_c^*, \bar{D}^*\Sigma_c, \bar{D}^*\Sigma_c^*$.
- $S - D$ mixing induced by the Tensor force (S_{12})

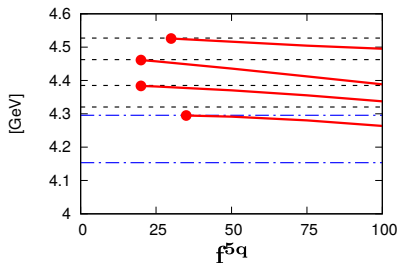
Result: V^{5q} (i) + (ii) + (iii)

- Hidden-charm: $V^{5q} = V_{(i)}^{5q} + V_{(ii)}^{5q} + V_{(iii)}^{5q}$

$$J^P = 1/2^-$$



$$J^P = 3/2^-$$



Volume integrals of the potentials

- Bound and Resonant states appears for $f^{5q} \gtrsim 25$
⇔ Large? Small?

Volume integrals of the potentials

- Bound and Resonant states appears for $f^{5q} \gtrsim 25$
 \Leftrightarrow Large? Small?

▶ Volume integral $V(q = 0) = \int V(r) dr^3$

Comparison with the NN interaction (Bonn potential)

R. Machleidt, K. Holinde and C. Elster, Phys. Rept. **149**, 1 (1987).

$$\left| V_{f=25}^{5q}(0) \right| = 1.1 \times 10^{-4} \text{ MeV} \sim 0.03 |C_{NN}^{\sigma}(0)|$$

(C_{NN}^{σ} : Central force of σ exchange)

- $\left| V_{f=25}^{5q}(0) \right|$ is **much smaller** than $|C_{NN}^{\sigma}(0)|$.

However, the bound and resonant states are obtained!